
Universal Blind Quantum Computation

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with

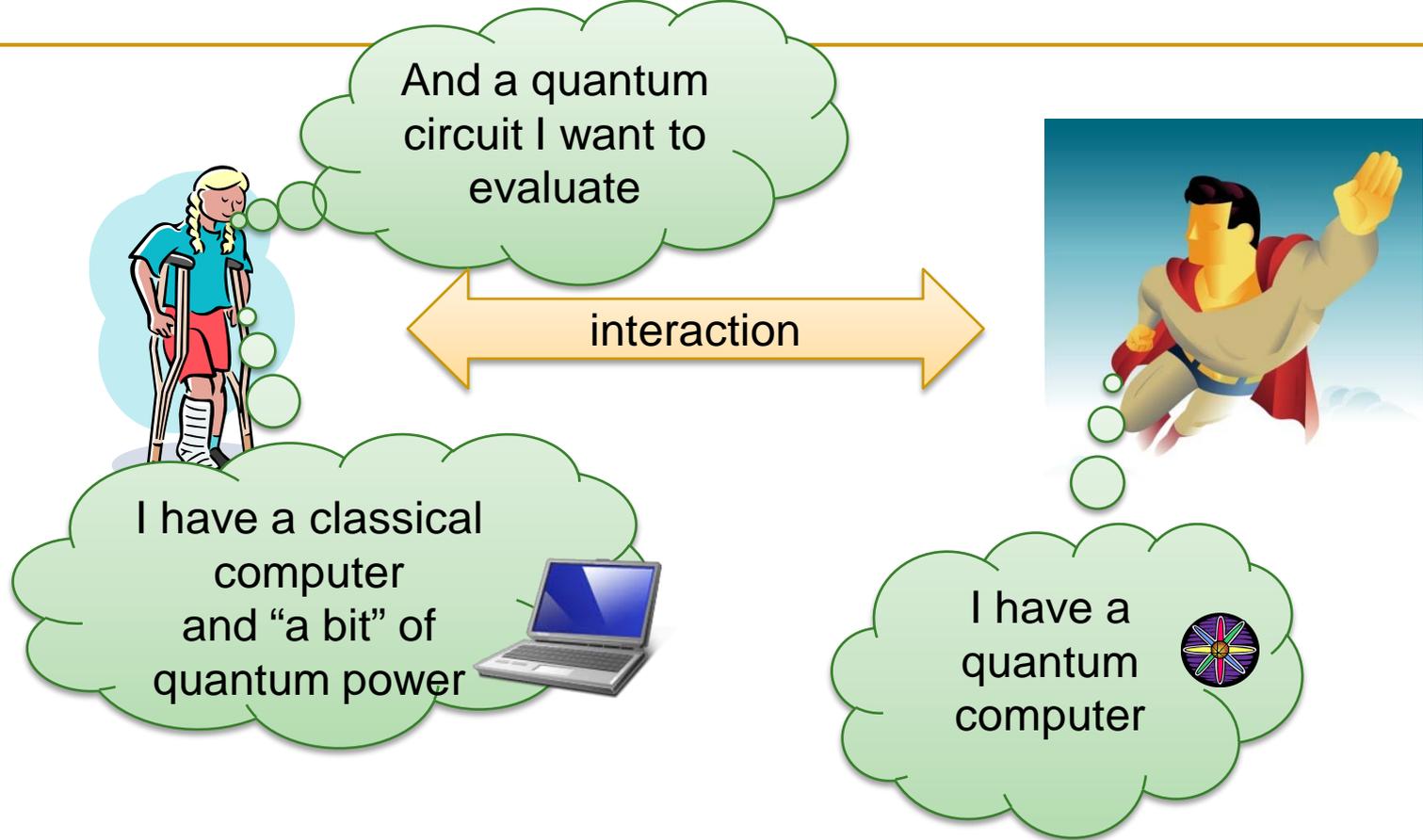
Joseph Fitzsimons (Oxford)

Elham Kashefi (Edinburgh)

20??



- The year is 20??. A few centers around the world have managed to build quantum computers.
- They allow users to have remote access to their quantum computers.



Interactive Proofs

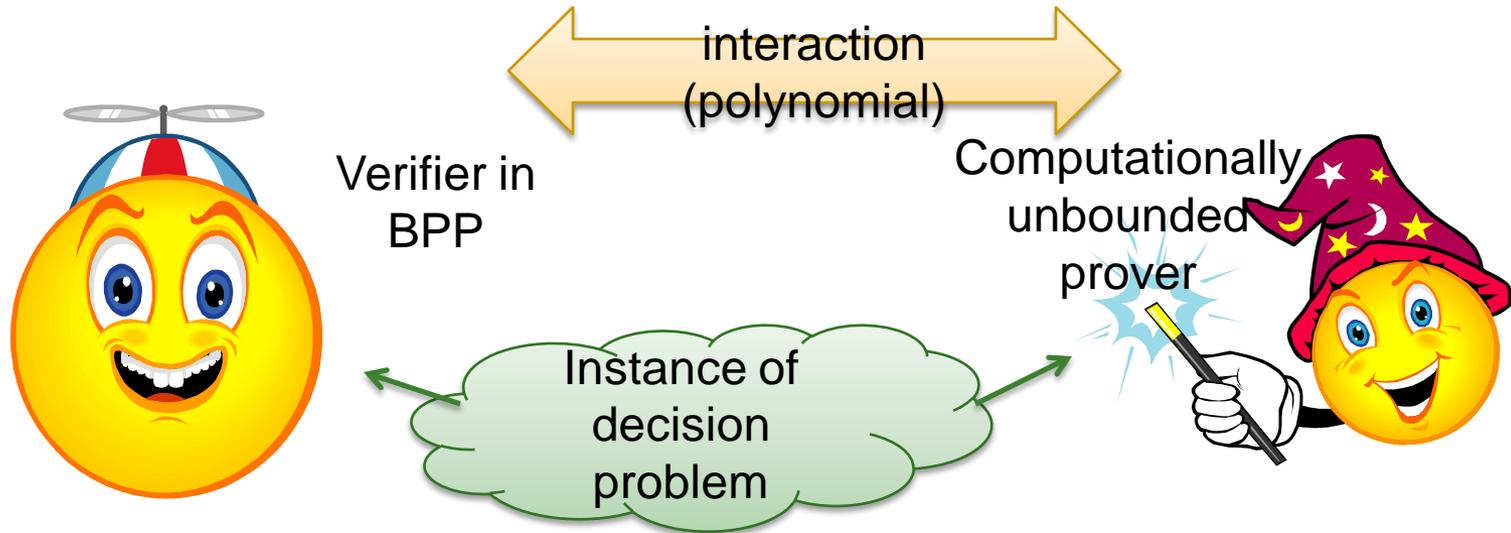
Cryptography

- How can Alice be convinced that the output provided by the quantum computer is correct?
- Can she do this while keeping her input private?

Interactive proofs

...how useful is a cheating oracle?

Classical interactive proofs (IP)

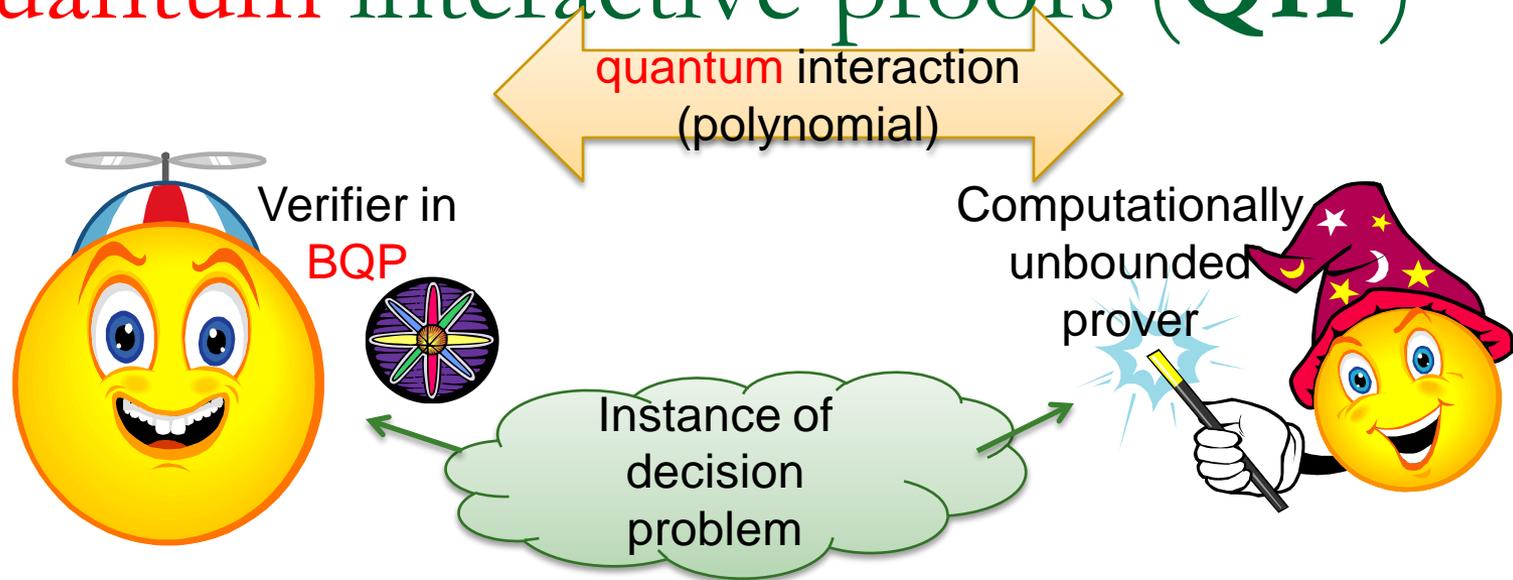


A language L is in **IP** if there exists a verifier such that:

- If the answer is "yes", the prover must be able to behave in such a way that the verifier accepts with probability at least $2/3$
- If the answer is "no", then however the prover behaves, the verifier must reject with probability at least $2/3$.

IP = PSPACE (Shamir, Lund-Fortnow-Karloff-Nisan 1990)

Quantum interactive proofs (QIP)



A language L is in **QIP** if there exists a verifier such that:

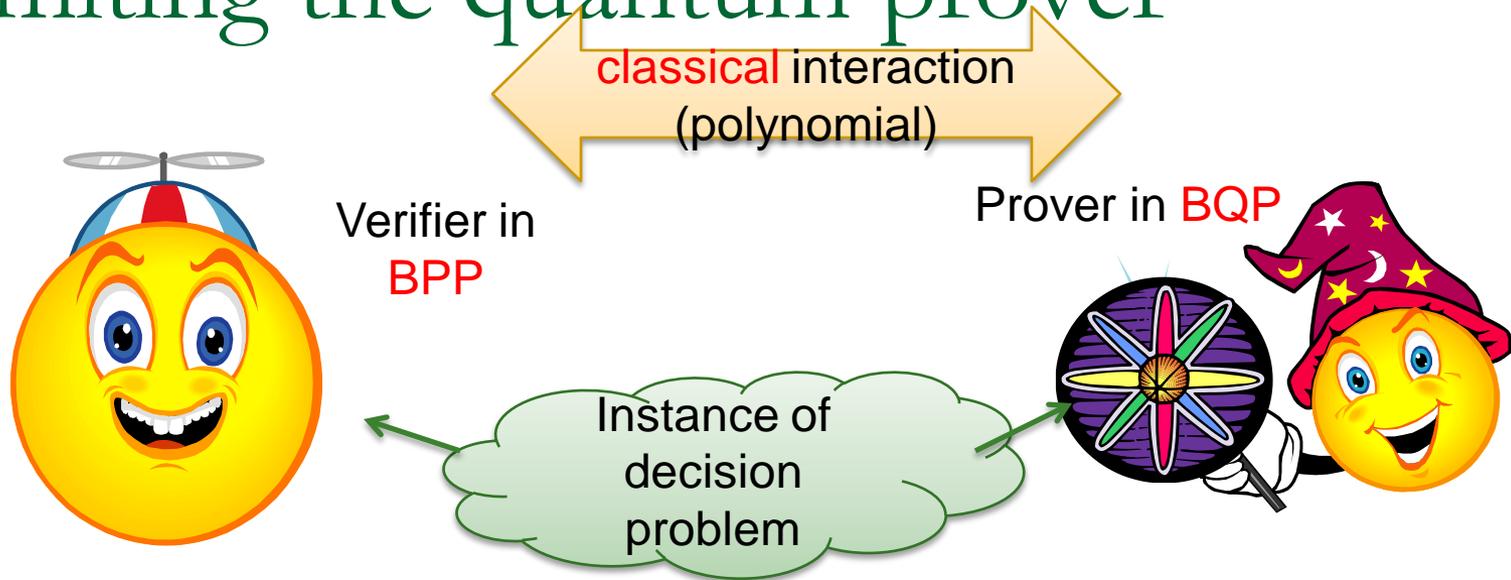
- If the answer is "yes," the prover must be able to behave in such a way that the verifier accepts with probability at least $2/3$
- If the answer is "no," then however the prover behaves the verifier must reject with probability at least $2/3$.

• **PSPACE** is in **QIP[3]** (Watrous 1999)

• **QIP[k] = QIP[3] = QIP** ($k \geq 3$) (Kitaev-Watrous 2000).

• Open question: Does **QIP** strictly contain **IP** (i.e. does quantum computation add any power to interactive proofs?)

Limiting the quantum prover



- Open question: what is the power of this type of scenario?

$$\text{IP}_{\text{BQP}} \stackrel{?}{=} \text{BQP}$$

- Our contribution: we give solutions to closely related problems:

1. Almost-classical verifier (has the additional power of generating random qubits from a fixed finite set):

$$\text{IP}_{\text{BQP}}^{|\theta\rangle} = \text{BQP}$$

Major open problem:
characterize the power of MIP^* .

2. Purely classical verifier, with **two** BQP provers that cannot communicate but that share entanglement

$$\text{MIP}_{\text{BQP}}^* = \text{BQP}$$

Cryptography

...what can be accomplished in the presence of an adversary?

Cryptography

- Quantum key distribution (QKD) (Bennett-Brassard 1984)
- Impossibility of Bit Commitment (Mayers, Lo-Chau 1995)
- Private Quantum Channels (Ambainis-Mosca-Tapp-de Wolf 2000)
- Quantum Authentication (Barnum-Crépeau-Gottesman-Smith-Tapp 2002)
- Multi-party computation (Ben-Or-Crépeau-Gottesman-Hassidim-Smith 2006)
- Cryptography in the bounded quantum-storage model (Damgård-Fehr-Salvail-Schaffner 2005)

Blind Quantum Computing



I have a classical computer and very limited quantum power



I have a quantum computer

**Our protocol achieves perfect privacy & detection of interfering Bob;
It can also be used for quantum inputs or outputs**

Motivations

■ Factoring

- Using Shor's algorithm, Alice can use Bob to help her factor an integer corresponding to an RSA public key
 - Bob won't learn whose private key he is breaking; in fact he won't even know that he is helping Alice factor.

■ BQP-Complete problem

- No known efficient method to verify solution: we therefore give a method to authenticate Bob's computation.

■ Processing quantum information

- Blind state preparation, blind measurement...

Previous work

Blind quantum computation

quant-ph/0309152

Pablo Arrighi^{1,*} and Louis Salvail^{2,†}

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- Publicly-known classical random-verifiable function
- Alice needs to be able to prepare and measure multi-qubit states
- Provides only *cheat sensitivity*

Previous work

arXiv:quant-ph/0111046

MIT-CTP #3211

Secure assisted quantum computation

Andrew M. Childs*
Center for Theoretical Physics
Massachusetts Institute of Technology
Cambridge, MA 02139, USA
(7 November 2001)

- Alice needs a quantum memory, and the ability to perform Pauli gates $X = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}, Z = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}$
 - Idea: she sends encrypted qubits to Bob who applies a known gate. Alice can decrypt the qubits while preserving the action of the gate. Repeat, cycling through universal set of gates.
- $$H = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 & 1 \\ 1 & -1 \end{pmatrix}, \pi/8 = \begin{pmatrix} 1 & 0 \\ 0 & \sqrt{i} \end{pmatrix}, CNOT = \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 0 & 1 \\ 0 & 0 & 1 & 0 \end{pmatrix}$$

Interactive Proofs For Quantum Computations

Dorit Aharonov*

Michael Ben-Or*

Elad Eban*

October 29, 2008

- Interactive proof with BQP prover, and nearly-classical verifier.
 - Verifier has a constant-size quantum computer
 - Protocol is also *blind*.

Our solution



Blind protocols that show:

$$\text{BQP} = \text{IP}_{\text{BQP}}^{|\theta\rangle}$$

$$\text{BQP} = \text{MIP}_{\text{BQP}}^*$$

High-level protocol



Classical input,
classical output

Input built
into circuit



- prepares qubits randomly chosen in

$$\left\{ \frac{1}{\sqrt{2}}(|0\rangle + e^{i\theta}|1\rangle) \mid \theta \in \left\{ \frac{n\pi}{4}, n = 0, 1, \dots, 7 \right\} \right\}$$

$|\uparrow\rangle$ $|\downarrow\rangle$ $|\leftarrow\rangle$ $|\searrow\rangle$
 $|\uparrow\rangle$ $|\rightarrow\rangle$ $|\swarrow\rangle$ $|\leftarrow\rangle$
 $|\nearrow\rangle$ $|\uparrow\rangle$ $|\uparrow\rangle$ $|\rightarrow\rangle$

repeat {

■ Classical computation



- Applies quantum operations and measurements

- Alice gets the output

Our technique

- Derived from Measurement Based quantum computing (MBQC)
[Raussendorf and Briegel, 2001]
- First time that a new functionality is achieved in MBQC.

The MBQC paradigm

Qubits are measured layer-by-layer...

How to convert any quantum circuit to MBQC:

1. Start with *cluster state*
2. Perform $\{|0\rangle, |1\rangle\}$ -basis measurements, depending on position of CNOT gates in quantum circuit
3. Perform x-y plane measurements adaptively, layer by layer



Final layer is output

Each qubit j has a target measurement angle ϕ_j

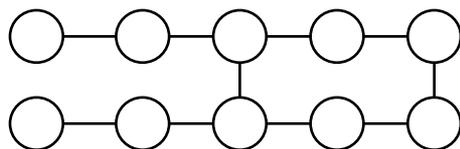
Each edge a two-qubit interaction

Measure in basis $\{\frac{1}{\sqrt{2}}(|0\rangle + e^{i\phi_j}|1\rangle), \frac{1}{\sqrt{2}}(|0\rangle - e^{i\phi_j}|1\rangle)\}$

ada $C - Z = \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & -1 \end{pmatrix}$
 $\phi'_j = (-1)^{s_x^j} \phi_j + \pi s_z^j$
 $(s_x^j \in \{0, 1\} \text{ and } s_z^j \in \{0, 1\})$
 depend on previous measurement outcomes

Getting rid of $\{|0\rangle, |1\rangle\}$ -basis measurements

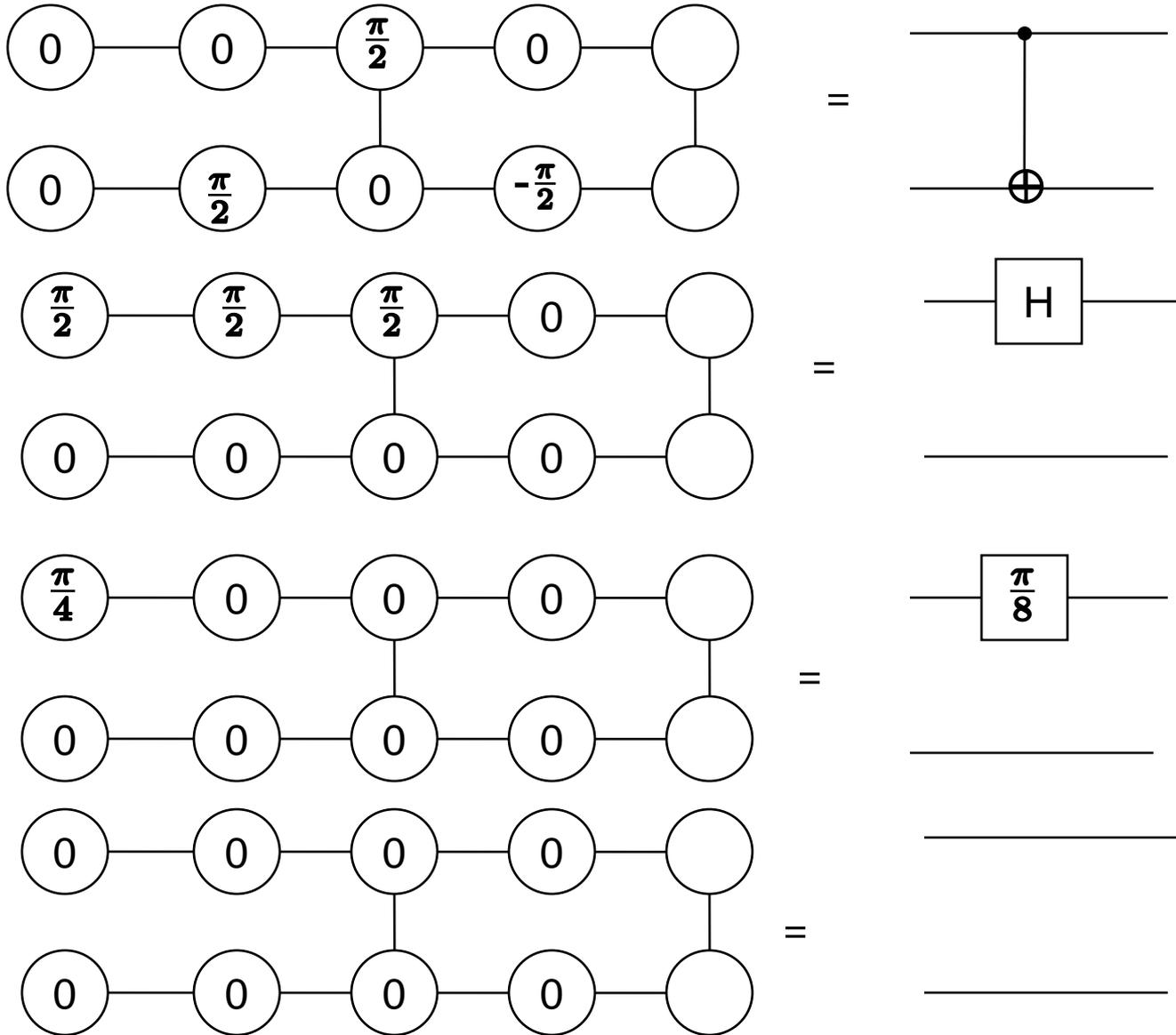
- We want to get rid of computational basis measurements that reveal the structure of underlying circuit
- We'll show that



yields universal set of gates: CNOT, H, and $\pi/8$

- Tiling the 2-qubit gate allows multiple inputs and multiple gates

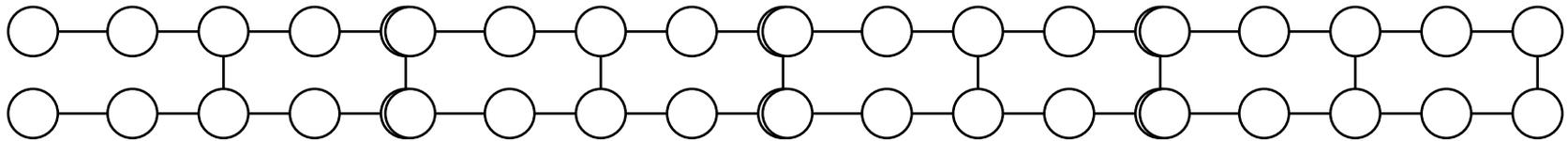
Getting rid of $\{|0\rangle, |1\rangle\}$ -basis measurements



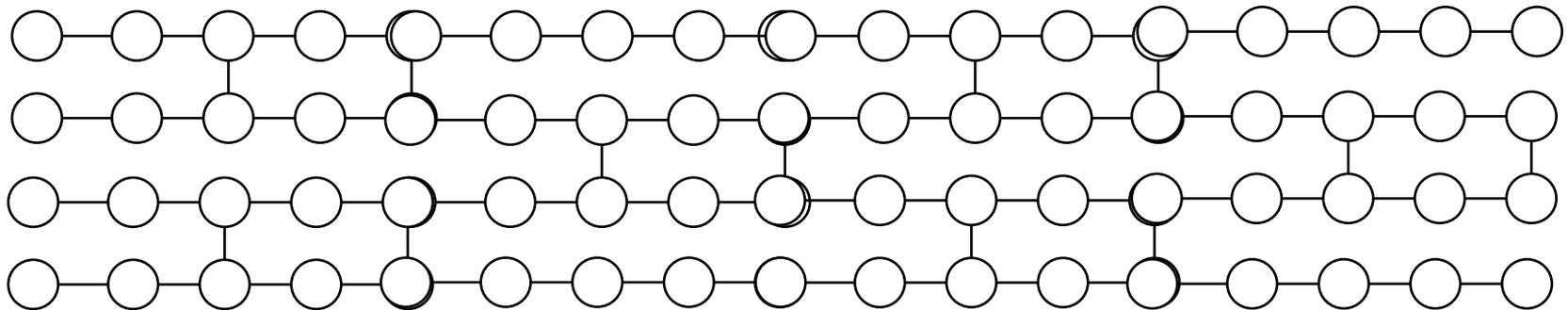
Getting rid of $\{|0\rangle, |1\rangle\}$ -basis measurements

The *brickwork* states

2-qubit circuit



4-qubit circuit



n -qubit circuit...

All measurements are integer multiples of $\frac{\pi}{4}$.

Blin... ol



Alice's Z-rotation..

...commutes with Bob's control-Z.

- prepares qubits randomly chosen in

$$\{\frac{1}{\sqrt{2}}(|0\rangle + e^{i\theta}|1\rangle) \mid \theta \in \{\frac{n\pi}{4}, n = 0, 1, \dots, 7\}\}$$

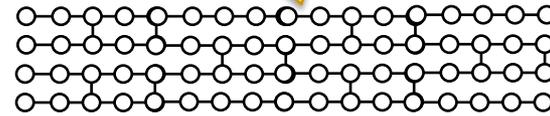
$$\begin{aligned} &|\uparrow\rangle \quad |\downarrow\rangle \quad |\leftarrow\rangle \quad |\searrow\rangle \\ &|\uparrow\rangle \quad |\rightarrow\rangle \quad |\leftarrow\rangle \quad |\swarrow\rangle \\ &|\nearrow\rangle \quad |\uparrow\rangle \quad |\uparrow\rangle \quad |\rightarrow\rangle \end{aligned}$$

- chooses x-y plane measurement angles, adaptively, layer by layer

$$\delta = \phi' + \theta + \pi r$$

$$\phi' = (-1)^{s_x} \phi + \pi s_z$$

- entangling according to brickwork state

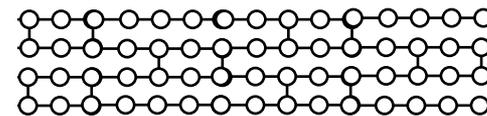


Measuring in δ basis cancels out Z-rotation

$$\delta_1, \delta_2, \delta_3, \delta_4$$

- single-qubit measurements in basis

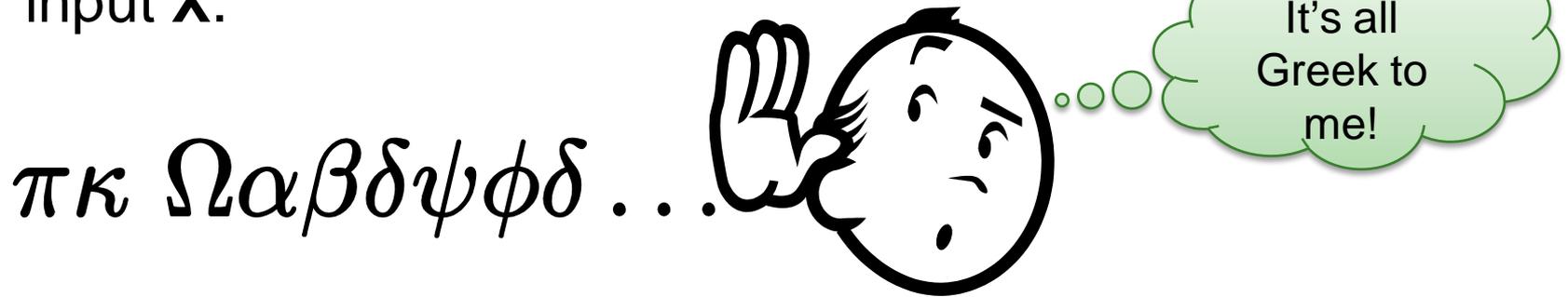
$$\{\frac{1}{\sqrt{2}}(|0\rangle + e^{i\delta}|1\rangle), \frac{1}{\sqrt{2}}(|0\rangle - e^{i\delta}|1\rangle)\}$$



r random. r=1 flips Bob's measurement outcome. Alice can correct this.

Privacy

- Intuitively, we want that from Bob's point of view, all information received from Alice is independent of Alice's input \mathbf{X} .



- Bob does learn the dimensions of the brickwork state, giving an upper bound on the size of Alice's computation. He may also have some prior knowledge on \mathbf{X} .
- Hence, we need to prove that Bob's view of the protocol does not depend on \mathbf{X} , given his prior knowledge.

Privacy

- Formally:

We say that a protocol is *blind while leaking at most $L(\mathbf{X})$* if for any fixed $\mathbf{Y}=\mathbf{L}(\mathbf{X})$, the following two hold when given \mathbf{Y} :

1. The distribution of the classical information obtained by Bob is independent of \mathbf{X} .
2. The state of the quantum system obtained by Bob is fixed and independent both of \mathbf{X} and of the distribution of the classical information above.

- Theorem: Our protocol is blind, while leaking at most the dimensions of the brickwork state.

Privacy



- prepares qubits randomly chosen

$$\left\{ \frac{1}{\sqrt{2}}(|0\rangle + e^{i\theta}|1\rangle) \mid \theta \in \left\{ \frac{n\pi}{4}, n = 0, 1, \dots, 7 \right\} \right\}$$

Let $\theta' = \theta + \pi r$
and $\vec{\theta}' = (\theta'_1, \theta'_2, \theta'_3, \dots)$

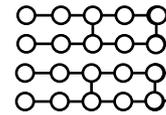
- chooses x-y plane measurement angles, adaptive layer by layer

$$\delta = \phi' + \theta + \pi r$$

$$\phi' = (-1)^{s_x} \phi + \pi s_z$$

Let **A** be the quantum system initially sent from Alice to Bob

entangled brickwork



Fix $\vec{\delta}$. Because **r**'s are random, for each qubit of **A**, one of the following two has occurred:

$$r = 0 \text{ so } |\psi\rangle = \frac{1}{\sqrt{2}}(|0\rangle + e^{i(\delta-\phi')})|1\rangle.$$

$$r = 1 \text{ so } |\psi\rangle = \frac{1}{\sqrt{2}}(|0\rangle - e^{i(\delta-\phi')})|1\rangle.$$

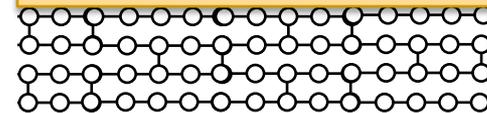
Hence when **r** is unknown, **A** consists of copies of the two-dimensional completely mixed state, which is fixed and independent of $\vec{\phi}$.

Let $\vec{\delta} = (\delta_1, \delta_2, \delta_3, \dots)$ be the classical information that Bob gets during the protocol

Hence $\vec{\delta} = \vec{\phi} + \vec{\theta}'$

$\vec{\theta}'$ is random, so $\vec{\delta}$ and $\vec{\phi}$ are independent

angle basis $\left\{ \frac{1}{\sqrt{2}} \right\}$



Detecting an interfering Bob

For classical outputs that cannot easily be verified

- ❑ Double the number of wires, randomly adding $N/2$ wires in $|0\rangle$ and $N/2$ wires in $|1\rangle$.
- ❑ An actively interfering Bob is caught with probability at least $\frac{1}{2}$. Repeat s times.
- We also have a fault-tolerant version that additionally provides authentication for quantum inputs and outputs.

Interactive proof



Verifier in **BPP** +
random qubits



Prover in **BQP**



- The blind protocol is as an interactive proof for any problem in BQP.

It follows:

$$\text{BQP} \subseteq \text{IP}_{\text{BQP}}^{|\theta\rangle}$$

Trivially,

$$\text{BQP} \supseteq \text{IP}_{\text{BQP}}^{|\theta\rangle}$$

Hence,

$$\text{BQP} = \text{IP}_{\text{BQP}}^{|\theta\rangle}$$

Multi-prover interactive proofs



$\tilde{\theta}_1, \tilde{\theta}_2, \dots$

Cheating is detected by the authentication protocol

$$\left\{ \frac{1}{\sqrt{2}}(|0\rangle + e^{i\tilde{\theta}_j}|1\rangle), \frac{1}{\sqrt{2}}(|0\rangle - e^{i\tilde{\theta}_j}|1\rangle) \right\}$$

$$\frac{1}{\sqrt{2}}(|00\rangle + |11\rangle)^{\otimes N}$$



Classical part of blind QC using

$$\theta_j = \tilde{\theta}_j + m_j \pi$$



Our result:

$$\text{BQP} \subseteq \text{MIP}_{\text{BQP}}^*$$

Trivially,

$$\text{BQP} \supseteq \text{MIP}_{\text{BQP}}^*$$

Hence, $\text{BQP} = \text{MIP}_{\text{BQP}}^*$

Open questions

- Is quantum communication required for blind quantum computation?

- $IP_{BQP} \stackrel{?}{=} BQP$



Thank you